

# **DOING PHYSICS WITH PYTHON**

## **DYNAMICAL SYSTEMS [1D]**

### **INTRODUCTION**

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### **DOWNLOAD DIRECTORIES FOR PYTHON CODE**

**[Google drive](#)**

**[GitHub](#)**

**ds25L1.py** Damped driven mass-spring system

**Jason Bramburger**

History and Preliminaries - Dynamical Systems | Lecture 1

<https://www.youtube.com/watch?v=Gw19VCtHYcs>

### **INTRODUCTION**

As an introduction to dynamical system, an oscillating mass-spring system will be considered. A dynamical system can be modelled by a set of differential equations and the behaviour of the system is completely described by the solutions of the differential equations. The mass-spring system is an example of an initial value problem

since you need to know the equations and the initial conditions. The equations are expressed in a set of state variables  $x_1, x_2, x_3, \dots$

## Mass – Spring system

### Simple Harmonic Motion (SHM)

### Damped Harmonic Motion (DHM)

### Forced Harmonic Motion (FHM)

The governing equation for the oscillating mass-spring system (no external forcing) is

$$(1) \quad m \ddot{x} + b \dot{x} + k x = 0$$

where  $m$  is the mass of the oscillating object,  $b$  is the damping constant and  $k$  is the spring constant. The state variable  $x$  is the displacement from the equilibrium position ( $x = 0$ ) along the X-axis. The equilibrium (steady-state) position is the only fixed point where  $x_{ss} = 0$ . This is **autonomous** equation since it does not have any terms which are an explicit function of time  $t$ .

Equation 1 is a 2<sup>nd</sup> order ODE and to solve it, the equation is expressed as a set of two 1<sup>st</sup> order ODEs where  $x_1$  and  $x_2$  are the state variables

$$(2) \quad \begin{aligned} \dot{x}_1 &= x_2 \\ \dot{x}_2 &= -(b/m)x_2 - (k/m)x_1 \end{aligned}$$

and the initial conditions are  $x_1(0)$  and  $x_2(0)$ . The state variable  $x_1$  is the displacement  $x$  of the oscillating object from the Origin ( $x = 0$ ),  $x_2$  is velocity  $v$  of the object, and  $\dot{x}_2$  is the object's acceleration  $a$ .

The natural frequencies and period for simple harmonic motion (SHM)  $b = 0$  are

$$(3) \quad \omega_0 = \sqrt{\frac{k}{m}} \quad f_0 = \left(\frac{1}{2\pi}\right) \sqrt{\frac{k}{m}} \quad T = 2\pi \sqrt{\frac{m}{k}}$$

The ODEs can be solved very easily using the Python function **odeint**. The output of the Code is displayed graphically.

#### # ds25L1.py

```
import numpy as np
from scipy.integrate import odeint
import matplotlib.pyplot as plt
from scipy.signal import find_peaks
import time
from numpy import pi, sin, cos, linspace, sqrt

plt.close('all')

tStart = time.time()

### SOLVE ODE x
def lorenz(t, state):
    x1, x2 = state
    dx1 = x2
    dx2 = -(b/m)*x2 - (k/m)*x1
    return dx1, dx2

### SETUP
u0 = [0.1,0]
tMax = 10; N = 9999
m = 2
b = 0.5
k = 9
```

$$w = \sqrt{k/m}$$

$$T = 2\pi/w$$

$$f = 1/T$$

`### SOLVE ODE`

`t = linspace(0,tMax,N)`

`sol = odeint(lorenz, u0, t, tfirst=True)`

`x = sol[:,0]`

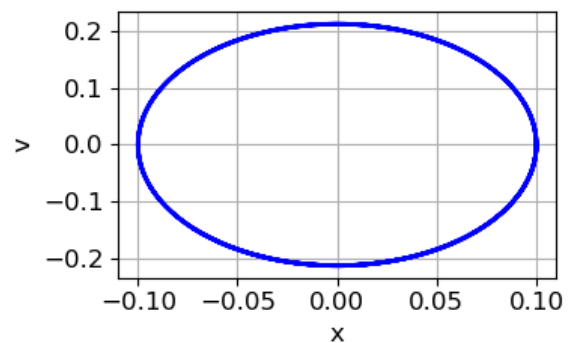
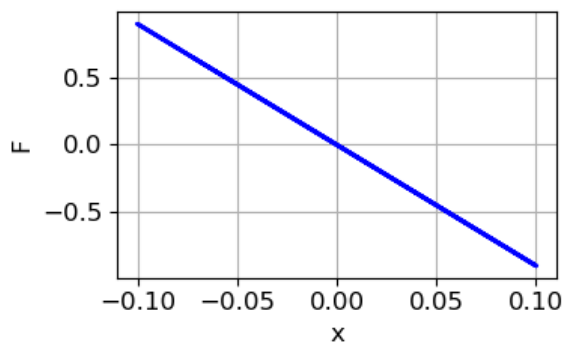
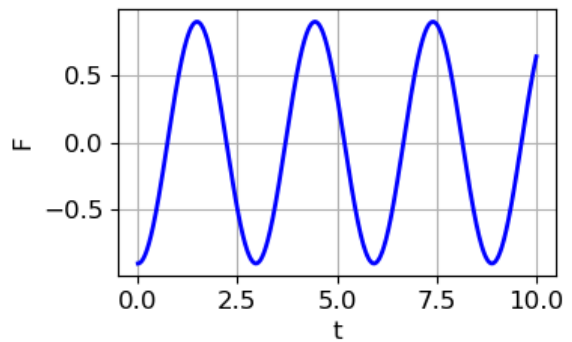
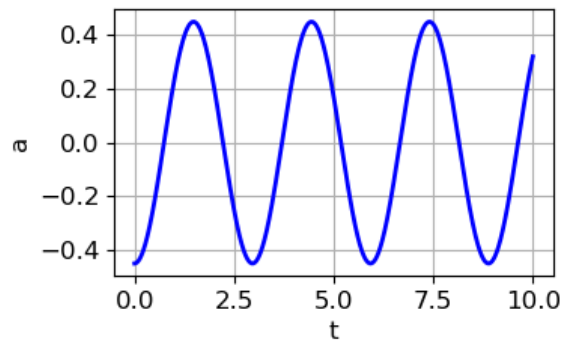
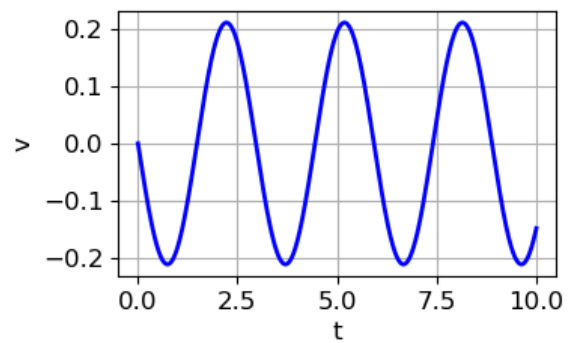
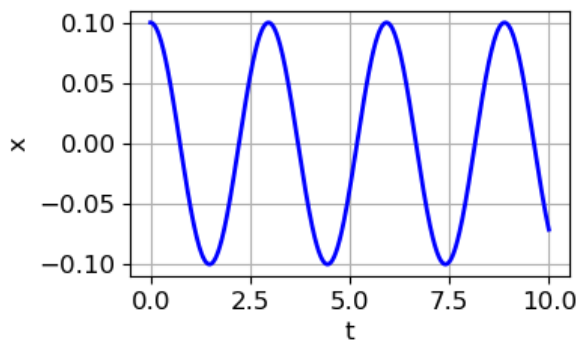
`v = sol[:,1]`

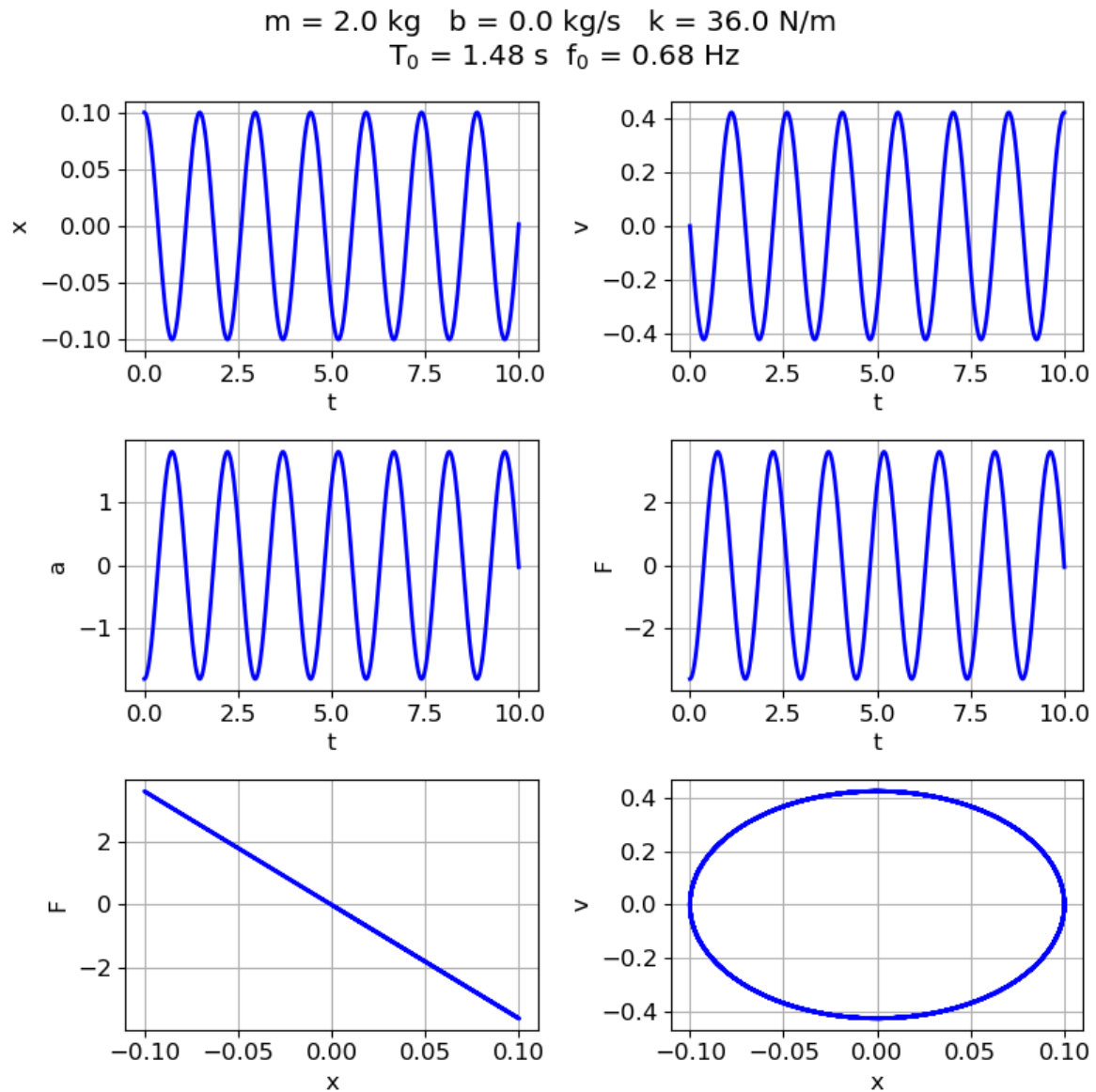
`a = -(b/m)*v - (k/m)*x`

`F = m*a`

$$m = 2.0 \text{ kg} \quad b = 0.0 \text{ kg/s} \quad k = 9.0 \text{ N/m}$$

$$T_0 = 2.96 \text{ s} \quad f_0 = 0.34 \text{ Hz}$$

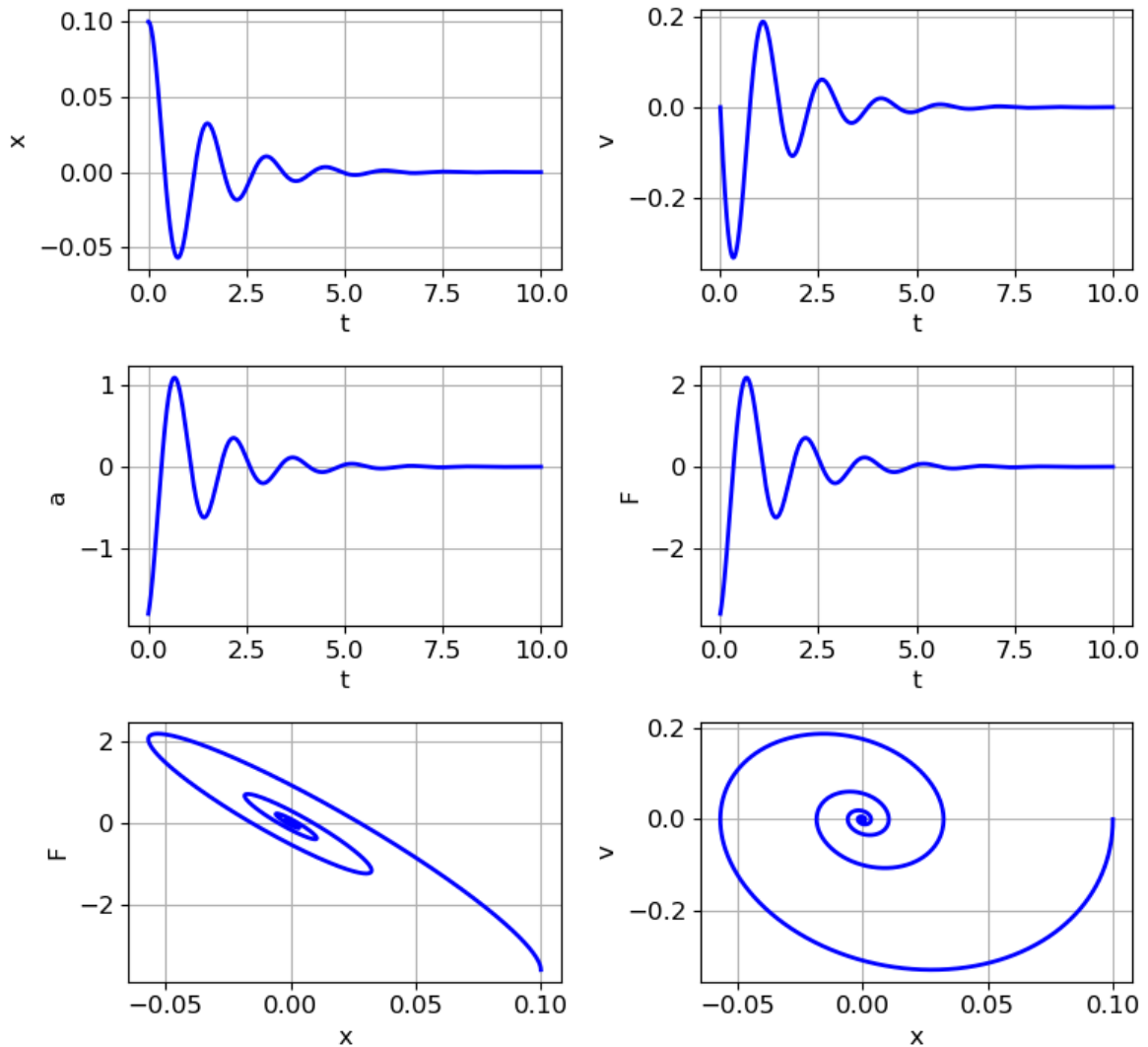




Increasing the value of  $k$  from 9 to 36 results in an increase in the frequency and a decrease in the period of oscillation.

$$m = 2.0 \text{ kg} \quad b = 3.0 \text{ kg/s} \quad k = 36.0 \text{ N/m}$$

$$T_0 = 1.48 \text{ s} \quad f_0 = 0.68 \text{ Hz}$$



Damping: The oscillations decay until the object comes to rest at the Origin  $x = 0$ .

By examining the  $x$  vs  $F$  graph, we see that the Origin acts as a basin of attraction because the force acting on the object is always pulling it towards the Origin.

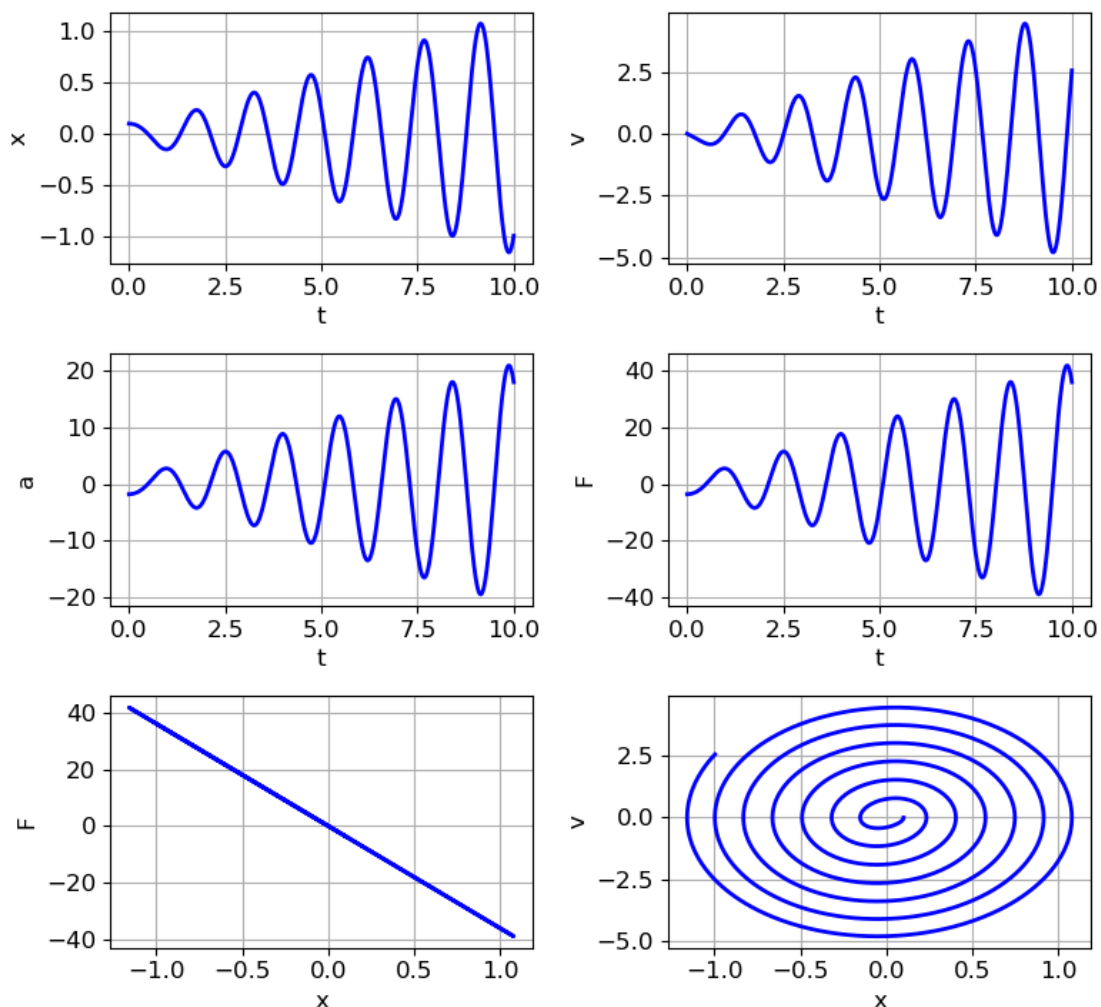
When an external sinusoidal force acts on the mass-spring system, then the equation of motion must include a time dependent term and the equation is now a **nonautonomous** equation

$$(3) \quad m \ddot{x} + b \dot{x} + k x = A_D \cos(\omega_D t)$$

where  $A_D$  is the strength of the driving force and the driving frequency is  $f_D$  ( $\omega_D = 2\pi f_D$ ). The set of 1<sup>st</sup> order ODEs are

$$(4) \quad \begin{aligned} \dot{x}_1 &= x_2 \\ \dot{x}_2 &= -(b/m)x_2 - (k/m)x_1 + A_D \cos(\omega t) \end{aligned}$$

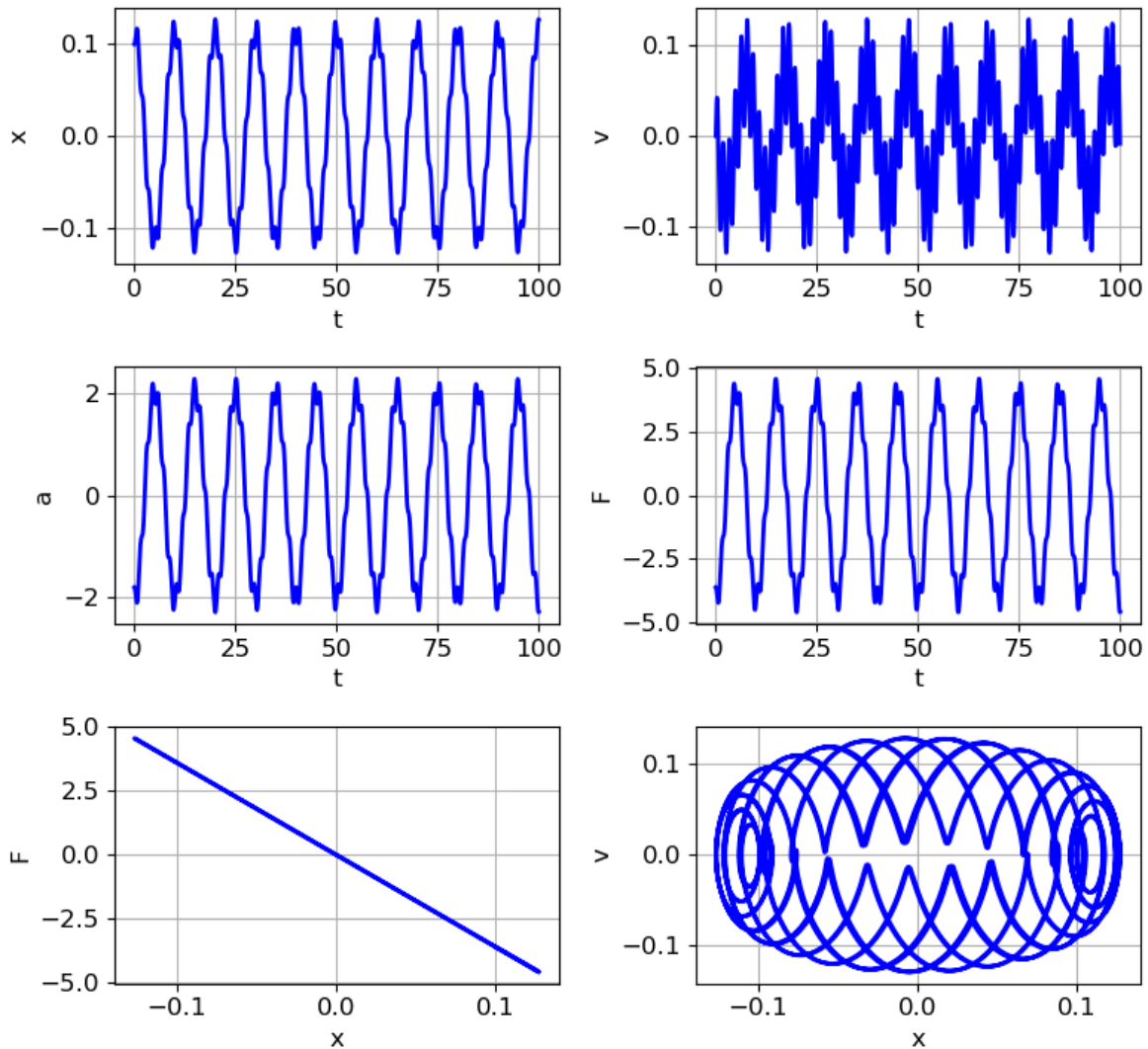
$$m = 2.0 \text{ kg} \quad b = 0.0 \text{ kg/s} \quad k = 36.0 \text{ N/m} \\ T_0 = 1.48 \text{ s} \quad f_0 = 0.68 \text{ Hz} \quad T_D = 1.45 \text{ s} \quad A_D = 1.00 \text{ m/s}$$



$b = 0 \quad f_D \approx f_0 \Rightarrow$  oscillations grow and grow (**resonance**)

$$m = 2.0 \text{ kg} \quad b = 0.0 \text{ kg/s} \quad k = 36.0 \text{ N/m}$$

$$T_0 = 1.48 \text{ s} \quad f_0 = 0.68 \text{ Hz} \quad T_D = 10.00 \text{ s} \quad A_D = 2.00 \text{ m/s}$$



When the driving frequency  $f_D$  is not near the natural frequency  $f_0$ , then oscillations do not grow and the system oscillates near the driving frequency.

**View**

[Chaotic dynamical systems: a damped driven pendulum](#)